The Mechanism and Kinetics of Growth of the Superconducting Compound Nb₃Sn

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A study has been made of the mechanism and kinetics of formation of $Nb₃Sn$ from the elemental components. The $Nb₃Sn$ forms partly by diffusion and partly by a solution/ deposition mechanism which depends on thermal gradient mass transfer. The effect of this is to modify the growth equation to $x = kt^{0.36}$ over the temperature range 950 to 1150° C. The temperature dependence of these two processes, given by the difference between the activation energies for diffusion and solution, is -9.7 kcal/g atom (-0.42 eV/atom) so that the thickness of the $Nb₃Sn$ layer produced in any given time decreases with increasing temperature.

Various experimental factors are discussed in terms of their influence on the rate of growth of the layer.

1. Introduction

During the last few years the compound $Nb₃Sn$ has become increasingly important as a highfield superconductor because of its high critical field and critical temperature. Unfortunately from a technological point of view, these good physical properties are accompanied by extreme brittleness. The only way of obtaining practically useful conductors which can be bent round reasonably small radii and which make use of the good superconducting properties of $Nb₃Sn$ is to produce them in a tape form containing a very thin layer of the compound close to the neutral axis. In this manner considerable flexibility can be retained without risk of damage to the superconducting layer.

In this connection, this paper presents a study of the formation of thin layers of $Nb₃Sn$ by direct reaction between niobium and liquid tin.

2. Experimental Techniques

A series of preliminary sessile-drop experiments was carried out to see if the wetting processes had any effect on the subsequent reaction. Small pieces of tin weighing 0.5 g were placed on 0.125 in. thick \times 0.5 in. diameter niobium discs in a tantalum resistance furnace and were heated to temperatures between 700 and 1200° C for periods of $\frac{1}{2}$ h to 8 h. A vacuum better than *202*

 5×10^{-5} torr was used up to 1100° C and an atmosphere of argon above this temperature.

The greater part of the work was concerned with the formation kinetics of $Nb₃Sn$ and was carried out on niobium-tin diffusion couples made by drilling a piece of 0.2 in diameter tin rod and inserting a piece of 0.125 in diameter niobium rod in the hole. Both pieces of metal were ultrasonically cleaned in trichloroethylene before being sealed off in silica capsules under a vacuum better than 10^{-5} torr. The specimens were heattreated in a horizontal resistance furnace for periods varying from a few minutes to 140 h.

For short time heat-treatments, e.g. 15 min, the time was considered to start when the furnace temperature returned to the original setting after the insertion of the specimen. The heating-up interval was 5 to 10 min, and must therefore be carefully considered when interpreting the structures observed in those specimens. Short-time experiments were also carried out using pressed niobium cups containing tin pellets in a high-frequency induction furnace. This technique provided sufficiently fast heating and cooling rates for realistic diffusion times of the order of one minute to be achieved.

After heat-treatment the couples were sectioned and polished using metallographic techniques developed in this laboratory for the preparation of niobium alloys [1]. The phases were distinguished by anodising in 10% citric acid at 22.5 V [2].

The equilibrium diagram on which this study is based was determined in this laboratory [3].

Throughout the work, the materials used were spectroscopically pure tin supplied by Johnson Matthey & Co, and a commercial grade of niobium, the analysis of which is given at the end of this paper.

3. Experimental Results

Examination of the drop-substrate interfaces after reaction did not show any irregular **effects** and it was concluded that the wetting processes presented no obstacle to the subsequent diffusion.

Different substrate surface preparations were used, namely, as ground, as polished, and as polished and anodised. No effect was observed on the rate of spreading or drop profile.

The growth kinetics of $Nb₃Sn$ in the range 950 ~ to 1200 ~ C were investigated using **the** diffusion couples. During the metallographic examination of the specimens it was noticed that the growth of the Nb₃Sn varied around the circumference of the rod, having in places a broken and crystalline appearance at the $Nb₃Sn/$ tin interface. Fig. 1 shows a layer of $Nb₃Sn$ of

Figure 1 Smooth layer of Nb_aSn formed after 4 h at 950° C $(X 610)$.

uniform thickness while on another area of the same rod, a rough layer such as that shown in fig. 2 was found. In general, the region of uniform thickness extended around the part of **the** circumference adjacent to the silica capsule. The layer on the remaining two-thirds or so of **the** circumference was rough. Because of the variation in thickness of the $Nb₃Sn$, it was not possible to follow the rate of formation by measuring **the** width of the layer in the normal manner. The anodising technique was therefore slightly modified to give a colour range which could be differentiated by a Metals Research "Quantitative Television Microscope". This instrument was used to measure the area of the $Nb₃Sn$ layer at eight equally-spaced points around **the** circumference and an average thickness of the NbaSn was calculated.

Figure 2 Irregular layer of Nb₃Sn formed after 4 h at 950° C $(x 610)$.

It is considered that the variation in growth of the $Nb₃Sn$ is due to a solution/deposition mechanism. The experimental method used will inevitably lead to thermal gradients within the capsule. Since the solubility of niobium in tin increases with temperature [4], the layer will be dissolved in the hot regions of the couple near **the** capsule wall where it lies on the furnace floor and the niobium will be deposited as $Nb₃Sn$ in the cooler regions. There will thus be thermal gradient mass transfer of niobium leading to an increased thickness of the layer in the cooler parts of the diffusion couple. Fig. 3 is a scanning electron microscope photograph of the surface of the Nb₃Sn after the overlying tin layer has been etched off. The surface of the $Nb₃Sn$ is seen to be crystalline in form, many of the crystals showing a large, well developed form which is characteristic of growth from the liquid. Further evidence of this mechanism is provided by **the** presence of small crystals of $NbSn₂$ in the tin 203

after solidification. Since $NbSn₂$ is not stable above 850° C [3], these must have been precipitated on cooling, confirming the presence of an appreciable solubility of niobium in the tin at 950° C.

Figure 3 Scanning electron micrograph of Nb₃Sn crystals $(X 2180)$.

If the growth of the $Nb₃Sn$ is controlled by diffusion then the thickness of the layer should follow a \sqrt{t} law, i.e.

$$
x = \mathbf{k}t^{0.5} \tag{1}
$$

where x is the layer thickness after a time t and k is a constant. Straight lines of the form

$$
\log x = m \log t + c \tag{2}
$$

were therefore fitted to the data by the method of least squares and table I shows the slopes and the intercepts at one hour for each temperature. The slope remains fairly constant at 0.34 to 0.39 over the range 950 \degree to 1150 \degree C. The results at 1200 \degree C do not fit this pattern, presumably because of the large amount of $Nb₃Sn$ deposited in the liquid tin at this temperature. Hence, from equation 2 the growth of the $Nb₃Sn$ layer between 950 $^{\circ}$ and 1150° may be written

$$
x = c' t^{0.36} \tag{3}
$$

where c' is a constant given by antilog c for any particular temperature and 0.36 is the average slope over the temperature range.

The index 0.36 in equation 3 indicates that the layer growth is slower than might be expected if theprocess were diffusion controlled(cf, equation 1). The constants in these two equations, k and 2O4

c', are the rate constants for the processes involved and contain the temperature dependence of the reaction rate.

4. Discussion

The effect of the solution/deposition mechanism on the growth kinetics of the $Nb₃Sn$ has been taken into account by considering a model system in which the solution and deposition regions are equal in area.

The rate of growth of the layer is given by

$$
\frac{\mathrm{d}x}{\mathrm{d}t} = \frac{\mathrm{k}_1}{x} - \mathrm{k}_2 \tag{4}
$$

for the region where solution is taking place and

$$
\frac{\mathrm{d}x}{\mathrm{d}t} = \frac{\mathrm{k}_1}{x} + \mathrm{k}_2 \tag{5}
$$

for the deposition region, k_1 describes the growth of the layer by diffusion, and k_2 describes the effect of the solution/deposition mechanism. It is reasonable to assume that k_2 is independent of x, but an approximation is being made in assuming that it is independent of t for all values of t . The solution rate will not become constant until a sufficient time has elapsed for the liquid tin to become saturated with niobium; however, for the purposes of the present treatment, this time has been neglected.

Integration and combination of equations 4 and 5 lead to an expression for t in terms of a polynomial function of \bar{x} , the average thickness, i.e.

$$
t = \frac{\bar{x}^2}{k_1} \left[\frac{5}{9} - \frac{28}{81} \frac{k_2 \bar{x}}{k_1} + \frac{34}{81} \frac{k_2^2 \bar{x}^2}{k_1^2} \ldots \ldots \right] \cdot
$$

Thus the effect of the solution/deposition mechanism has replaced the term x^2 in the usual diffusion equation $x^2 = 2Dt$ by a polynomial series. The index 0.36 in equation 3 is an experimental approximation to the series obtained by expressing x as a function of t .

From equation 4 it is inferred that there is a

maximum thickness of the solution region which will be reached when the diffusion growth is counterbalanced by solution.

Then $x_{\text{max}} = \frac{k_1}{k_2}$.

From the experimental data, a reasonable estimate of x_{max} at 950° C is 25 μ m, i.e. k₁/k₂ $= 2.5 \times 10^{-3}$ cm. Combining this value with the experimental results at 950 $^{\circ}$ C, k₁, the diffusion coefficient is evaluated as $\sim 10^{-10}$ $cm²/sec.$

A value of the observed resultant activation energy may be obtained from the slope of a graph of the rate of the process against $1/T^{\circ}$ K (fig. 4). If the layer thickness at 1 h is used as a measure of the rate of the process then the apparent activation energy is calculated to be -9.7 kcal/g atom. The negative value of the temperature dependence is an expression of the fact that as the temperature increases the solution part of the process increases more rapidly than the diffusion growth part, thus for any given time the average layer thickness decreases with rising temperature.

Figure 4 Temperature dependence of the growth rate of Nb₃Sn.

5. Factors affecting the growth of Nb_aSn 5.1. Temperature Gradients

The results of these experiments show clearly that a considerable amount of the $Nb₃Sn$ which forms does so by deposition from the liquid phase. The crystals then agglomerate and are slowly bonded together by the diffusion of niobium into the intermingling tin and its subsequent conversion to Nb₃Sn. It is not clear whether the niobium is transferred through the liquid tin by diffusion or by convection currents. A magnetic field of 3 kOe was applied during one

diffusion anneal in an attempt to damp out any convection currents in the liquid tin, but no difference was found in the mode of growth of the NbaSn. It has been calculated that the required temperature gradient to account for the necessary difference in thickness is about 10° at 950° C assuming that the transfer depends entirely on diffusion. This difference is larger than might have been expected in the small silica capsule inside the furnace.

If, however, the problem is considered in terms of diffusion through a static boundary layer of liquid followed by convective transfer, then the diffusion distance is much reduced since typical thicknesses of boundary layers vary from 10^{-1} cm to 10^{-3} cm. The mass flow may therefore be maintained by a much smaller concentration difference and hence the necessary temperature gradient is also reduced. For example for a static layer, 10^{-2} cm thick, the temperature difference is reduced to $\langle 1^{\circ}$ C.

The convective flow velocity required to transport the dissolved niobium is calculated to be of the order of 10^{-3} cm/sec. This low velocity explains the absence of any effect due to the magnetic field.

It is seen, therefore, that the solution/deposition mechanism requires very small driving forces, and will be difficult to eliminate in practice.

5.2. Amount and Composition of Liquid Tin

The saturation of the tin with niobium before reaction would have little effect, since it would allow the transport of the niobium to begin immediately without a preliminary delay while saturation was reached.

The volume of the tin is important because this influences the amount of niobium which can be held in solution. Temperature gradients will also be minimised in an experimental configuration which is physically small in size, thus helping to restrict the mass-transfer mechanism.

5.3. Surface Condition of the Niobium

It was concluded from the experiments on wetting that the surface preparation of the niobium was not important, but it has been found that the presence of a thickened oxide film has an effect upon the growth of the $Nb₃Sn$. Where the compound layer is smooth, the oxide film can frequently be seen adhering very close to the surface. When the niobium was heated in air for 24 h at 200° C before reaction, the

subsequent growth of the $Nb₃Sn$ was found to be dependent on the adhesion of the film. Fig. 5 shows that when the film breaks away, the growth is irregular. The position of the film is clearly marked by decoration with $NbSn₂$ crystals during cooling. The transition region demonstrates very clearly that a considerable amount of the niobium is consumed by solution in the tin.

Figure 5 Transition region between adhesion and nonadhesion of a surface oxide film (\times 610).

5.4. Heating Rates

The heating rates used for the diffusion couples are slow enough to allow the formation of appreciable amounts of $NbSn₂$ before the diffusion temperature is reached. After short diffusion times some $NbSn₂$ is still present together with the $Nb₃Sn$. The subsequent transformation of the $NbSn₂$ to $Nb₃Sn$ takes place in an irregular manner. In addition, some of the $NbSn₂$ also dissolves into the tin as the niobium solubility increases with increasing temperature. These two processes result in an irregular $Nb₃Sn/$ tin interface. Observations have shown that the formation of $Nb₃Sn$ is accomplished primarily by diffusion of the tin through the $Nb₃Sn$ into the niobium. Consequently any cracks or fissures in the $Nb₃Sn/tin$ interface cause indentations in the niobium/Nb3Sn interface and defects in the NbaSn layer do not close up. A faster heating rate would largely preclude the formation of $NbSn₂$ and would give a more uniform layer of NbaSn. Similarly a rapid cooling rate is also desirable, particularly if any excess tin is present 206

since this will react with the Nb₂S_n which has formed to give a decomposition product of NbSn₂.

5.5. Purity of Niobium

Niobium samples from different sources with purities varying from ordinary commercial grades to super-pure vacuum-annealed material were used in selected duplicate experiments to see if composition had any effect on formation of the NbaSn. No significant differences were observed.

6. Conclusions

(1) The growth of $Nb₃Sn$ by direct reaction between niobium and liquid tin is a complex reaction involving both diffusion and a solution/ deposition mechanism. The result of these two opposing processes has a negative temperature dependence which can be described by an apparent activation energy of $-$ 9.7 kcal/g atom. The detailed behaviour of the formation kinetics depends in practice upon the experimental configuration.

(2) The chemical diffusion coefficient for the formation of Nb₃Sn is $\sim 10^{-10}$ cm²/sec at 950° C.

(3) Because of the solution/deposition mechanism the volume of tin present is important. A minimum amount of tin leads to improved uniformity of the Nb₃Sn layers.

(4) Rapid heating rates are desirable to prevent the formation of $NbSn₂$ which subsequently transforms to an irregular $Nb₃Sn$ layer.

(5) Growth of the $Nb₃Sn$ layer occurs primarily at the niobium/Nb_aSn interface, suggesting that tin diffuses more rapidly than niobium through the $Nb₃Sn$ layer.

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Appendix

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